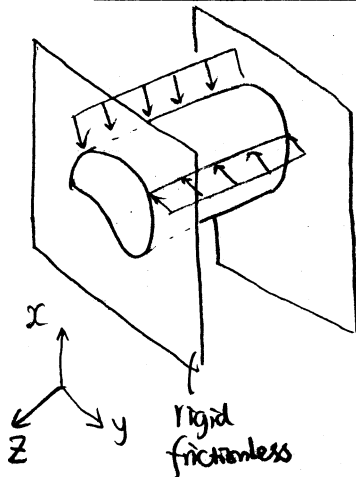


The equations of elasticity can be greatly simplified if we restrict the solution to be 2-Dimensional, i.e. does not depend on one coordinate (e.g. z).

There are several types of 2D elasticity problems:

- plane strain
- plane stress
- anti-plane strain (e.g. *straight screw dislocations)
- etc.

§1. Plane Strain



Consider a long bar with ends constrained by rigid-frictionless plates

Let the loads applied on the side of the bar be independent of z .

Then $u_z = 0$ every where
 u_x, u_y independent of z .

i.e. the unknowns of this problem is
 $u_x(x, y)$ and $u_y(x, y)$

We can easily show

$$\epsilon_{xz} = \frac{1}{2} \left(\frac{\partial u_x}{\partial z} + \frac{\partial u_z}{\partial x} \right) = 0, \quad \epsilon_{yz} = 0, \quad \epsilon_{zr} = 0$$

So the only non-zero strain components are

$\epsilon_{xx}, \epsilon_{yy}, \epsilon_{xy}$ — hence the name plane strain.

In isotropic elasticity, $\sigma_{xz} = \sigma_{yz} = 0$.

but in general $\sigma_{zz} \neq 0$

$$\epsilon_{zz} = -\frac{\nu}{E} \sigma_{xx} - \frac{\nu}{E} \sigma_{yy} + \frac{1}{E} \sigma_{zz} \rightarrow \sigma_{zz} = \nu(\sigma_{xx} + \sigma_{yy})$$

equilibrium condition

$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{yx}}{\partial y} + \frac{\partial \sigma_{zx}}{\partial z} + F_x = 0 \quad \left(\frac{\partial}{\partial z} = 0 \right)$$

$$\frac{\partial \sigma_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \sigma_{zy}}{\partial z} + F_y = 0$$

$$\frac{\partial \sigma_{xz}}{\partial x} + \frac{\partial \sigma_{yz}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z} + F_z = 0 \quad (\sigma_{xz} = \sigma_{yz} = 0, F_z = 0)$$

so the only non-trivial equations are

$$\boxed{\begin{aligned} \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{yx}}{\partial y} + F_x &= 0 \\ \frac{\partial \sigma_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + F_y &= 0 \end{aligned}}, \quad \sigma_{zz} = \nu(\sigma_{xx} + \sigma_{yy})$$

compatibility condition

$$\epsilon_{ij,kl} + \epsilon_{kl,ij} - \epsilon_{ik,jl} - \epsilon_{jl,ik} = 0$$

only one non-trivial equation when $i, j, k, l = x, y$

$$\boxed{\epsilon_{xx,yy} + \epsilon_{yy,xx} - 2\epsilon_{xy,xy} = 0}$$

Generalized Hooke's Law

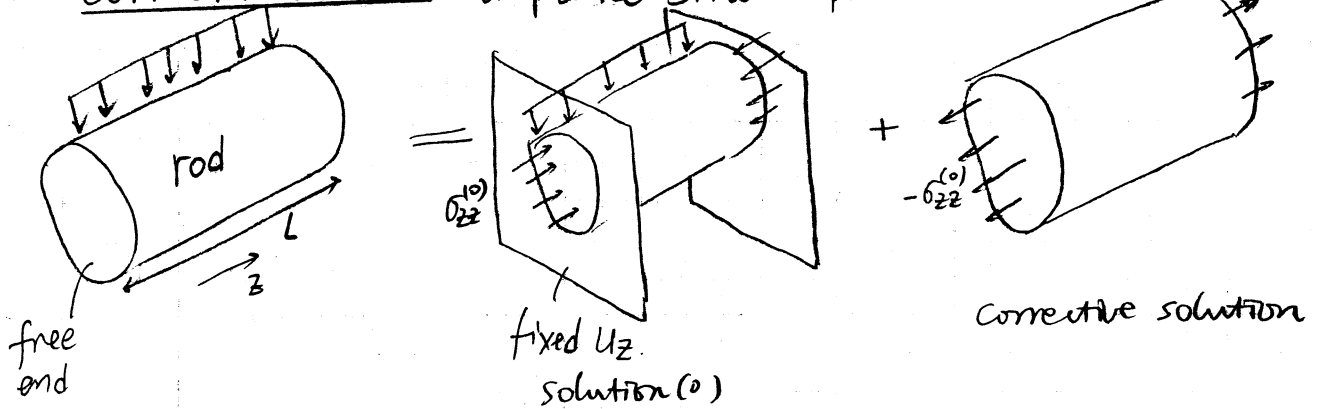
$$\epsilon_{xx} = \frac{1}{E} \sigma_{xx} - \frac{\nu}{E} \sigma_{yy} - \frac{\nu}{E} \sigma_{zz}$$

$$\epsilon_{yy} = -\frac{\nu}{E} \sigma_{xx} + \frac{1}{E} \sigma_{yy} - \frac{\nu}{E} \sigma_{zz}$$

$$\epsilon_{zz} = -\frac{\nu}{E} \sigma_{xx} - \frac{\nu}{E} \sigma_{yy} + \frac{1}{E} \sigma_{zz} = 0 \rightarrow \sigma_{zz} = \nu(\sigma_{xx} + \sigma_{yy})$$

$$\therefore \boxed{\begin{aligned} \epsilon_{xx} &= \frac{1-\nu^2}{E} \sigma_{xx} - \frac{\nu(1+\nu)}{E} \sigma_{yy} \\ \epsilon_{yy} &= -\frac{\nu(1+\nu)}{E} \sigma_{xx} + \frac{1-\nu^2}{E} \sigma_{yy} \end{aligned}}$$

* corrective solution to plane strains problem



In reality, we seldom have rods with both ends constrained by rigid, frictionless plates.

However, the plane strain solution is a good approximation to a long rod subjected to lateral loading that is z -independent.

To better describe the physical problem (rods with free end), a corrective solution should be added, in which axial forces are applied to the ends.

Because there may be a net force F at the end, the corrective solution is not necessarily limited to the neighborhood of the ends.

$$F = - \int_A \sigma_{zz}^{(0)} dA.$$

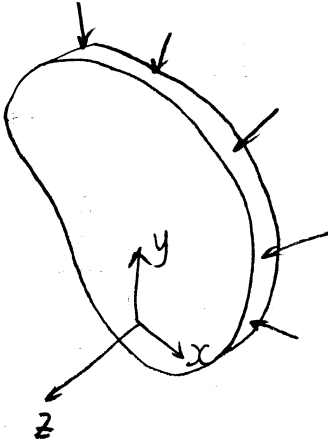
The corrective solution can be approximated by a simple solution

$$\sigma_{zz}^{(1)} = \frac{F}{A} = - \frac{1}{A} \int_A \sigma_{zz}^{(0)} dA.$$

$$\text{The rod length will change by } \Delta L = \frac{FL}{EA} = - \frac{L}{EA} \int_A \sigma_{zz}^{(0)} dA$$

* Of course, the "true" corrective solution is going to be more complicated by this. So we can imagine a correction solution to this correction solution, which should then be localized near the ends.

§2. Plane Stress



The opposite limit of plane strain problem (thick rods) is a thin film — plane stress problem.

Because $\sigma_{xz} = \sigma_{yz} = \sigma_{zz} = 0$ on the surface, if the film is sufficiently thin, we can expect

$$\sigma_{xz} \approx 0, \quad \sigma_{yz} \approx 0, \quad \sigma_{zz} \approx 0$$

everywhere inside the film.

Plane stress condition:

$$\sigma_{xz} = \sigma_{yz} = \sigma_{zz} = 0, \quad \frac{\partial}{\partial z} = 0.$$

Generalized Hooke's Law:

$$\epsilon_{xx} = \frac{1}{E} (\sigma_{xx} - \nu \sigma_{yy})$$

$$\epsilon_{xz} = \epsilon_{yz} = 0$$

$$\epsilon_{yy} = \frac{1}{E} (-\nu \sigma_{xx} + \sigma_{yy})$$

$$\epsilon_{zz} = -\frac{\nu}{E} (\sigma_{xx} + \sigma_{yy}), \quad \epsilon_{zz} \neq 0.$$

equilibrium condition:

$$\sigma_{xx,x} + \sigma_{yx,y} + F_x = 0$$

$$\sigma_{xy,x} + \sigma_{yy,y} + F_y = 0$$

Same as plane strain

compatibility condition:

$$\epsilon_{ij,kl} + \epsilon_{kl,ij} - \epsilon_{ik,jl} - \epsilon_{jl,ik} = 0$$

$$\epsilon_{xx,yy} + \epsilon_{yy,xx} - 2\epsilon_{xy,xy} = 0$$

$i=j=x, k=l=y$
Same as plane strain

$$* \begin{cases} \epsilon_{zz,xx} = 0 \\ \epsilon_{zz,yy} = 0 \\ \epsilon_{zz,xy} = 0 \end{cases} \quad \text{(cannot be satisfied all at the same time)}$$

$$\begin{cases} i=j=z, k=l=x \\ i=j=z, k=l=y \\ i=j=z, k=x, l=y \end{cases}$$

Because $\epsilon_{zz}(x,y) \neq 0$, we have more compatibility conditions.

§3. Equivalence between Plane Strain and Plane Stress

Both Plane strain and Plane Stress Problems seek solutions for $\sigma_{xx}, \sigma_{xy}, \sigma_{yy}$ and $\epsilon_{xx}, \epsilon_{xy}, \epsilon_{yy}$ as functions of x and y .

They satisfy the same equilibrium and compatibility conditions.

The only difference is in the Generalized Hooke's Law.

Plane Strain

$$\epsilon_{xx} = \frac{1-\nu^2}{E} \sigma_{xx} - \frac{\nu(1+\nu)}{E} \sigma_{yy}$$

$$\epsilon_{yy} = -\frac{\nu(1+\nu)}{E} \sigma_{xx} + \frac{1-\nu^2}{E} \sigma_{yy}$$

$$\epsilon_{xy} = \frac{1}{2\mu} \sigma_{xy}$$

Plane Stress

$$\epsilon_{xx} = \frac{1}{E} \sigma_{xx} - \frac{\nu}{E} \sigma_{yy}$$

$$\epsilon_{yy} = -\frac{\nu}{E} \sigma_{xx} + \frac{1}{E} \sigma_{yy}$$

$$\epsilon_{xy} = \frac{1}{2\mu} \sigma_{xy}$$

introduce Kolosov's constant κ

$$\boxed{\kappa = 3-4\nu}$$

$$\boxed{\kappa = \frac{3-\nu}{1+\nu}}$$

Then both Plane Stress and Plane Strain Problems satisfy the same condition

$$\epsilon_{xx} = \frac{1-\nu^2}{E} \sigma_{xx} - \frac{\nu(1+\nu)}{E} \sigma_{yy}$$

$$\epsilon_{yy} = -\frac{\nu(1+\nu)}{E} \sigma_{xx} + \frac{1-\nu^2}{E} \sigma_{yy}$$

Hence we can solve each 2D problem in either plane strain or in plane stress condition and find the other solution by replacing the Kolosov's constant.

§4. Airy Stress function

- Introduce a scalar function $\phi(x, y)$, and write different components of the stress tensor as derivatives of ϕ .
- We need to use 2nd derivatives, so that they transform as rank-2 tensor.
- Assume zero body force: $F_x = F_y = 0$
then the equilibrium condition becomes

$$\sigma_{xx,x} + \sigma_{xy,y} = 0 \quad \text{and} \quad \sigma_{yx,x} + \sigma_{yy,y} = 0$$

- These can be automatically satisfied by defining

$$\sigma_{xx} = \phi_{,yy}$$

$$\sigma_{yy} = \phi_{,xx}$$

$$\sigma_{xy} = -\phi_{,xy}$$

↑
notice the minus sign!

The equilibrium condition can be easily verified:

$$\begin{cases} \sigma_{xx,x} + \sigma_{xy,y} = \phi_{,yxx} - \phi_{,xyy} = 0 \\ \sigma_{yx,x} + \sigma_{yy,y} = -\phi_{,xyx} + \phi_{,xxy} = 0 \end{cases}$$

- So the only PDE we need to worry about is the compatibility condition.

$$\epsilon_{xx,yy} + \epsilon_{yy,xx} - 2\epsilon_{xy,xy} = 0$$

$$\begin{cases} \epsilon_{xx} = \frac{\kappa+1}{8\mu} \phi_{,yy} - \frac{3-\kappa}{8\mu} \phi_{,xx} \\ \epsilon_{yy} = -\frac{3-\kappa}{8\mu} \phi_{,yy} + \frac{\kappa+1}{8\mu} \phi_{,xx} \\ \epsilon_{xy} = -\frac{1}{2\mu} \phi_{,xy} \end{cases}$$

$$\frac{\kappa+1}{8\mu} (\phi_{,yyyy} + \phi_{,xxxx}) - \frac{3-\kappa}{8\mu} \cdot 2\phi_{,xxyy} + \frac{2}{2\mu} \phi_{,xyxy} = 0$$

$$\frac{\kappa+1}{8\mu} (\phi_{,xxxx} + \phi_{,yyyy} + 2\phi_{,xxyy}) = 0$$

$$\therefore \boxed{\nabla^4 \phi = 0} \quad \text{or} \quad \nabla^2 \nabla^2 \phi = 0 \quad \text{or} \quad \phi_{,xxxx} + \phi_{,yyyy} + 2\phi_{,xxyy} = 0$$

§5. Airy stress function in the presence of body forces.

When the body force can be written as derivatives of a potential

$$F_x = -\frac{\partial V}{\partial x} \quad F_y = -\frac{\partial V}{\partial y} \quad (\text{e.g. } V = \rho g y \text{ for gravity})$$

Then the Airy stress function can be defined as

$$\sigma_{xx} = \phi_{,yy} + V \quad \sigma_{yy} = \phi_{,xx} + V \quad \sigma_{xy} = -\phi_{,xy}$$

We can verify that the equilibrium condition

$$\sigma_{xx,x} + \sigma_{xy,y} + F_x = 0$$

$$\sigma_{yx,x} + \sigma_{yy,y} + F_y = 0$$

is automatically satisfied.

The compatibility condition becomes

$$\nabla^4 \phi = -2 \cdot \frac{\kappa-1}{\kappa+1} \cdot \nabla^2 V$$

in plane strain

$$\nabla^4 \phi = -\left(\frac{1-2\nu}{1-\nu}\right) \nabla^2 V$$

in plane stress

?

§6. Example 1

when there is no body force, we have $\nabla^4 \phi = 0$

Let's pick a trial solution

$$\phi(x,y) = \alpha x + \beta y + \gamma$$

It obviously satisfy the condition $\nabla^4 \phi = 0$.

What does this solution mean?

Let's find out the stress field.

$$\sigma_{xx} = \phi_{,yy} = ?$$

$$\sigma_{yy} = \phi_{,xx} =$$

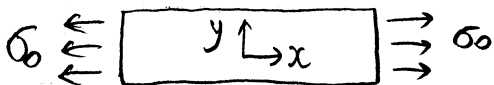
$$\sigma_{xy} = -\phi_{,xy} =$$

§7. Example 2.

Let's go to higher order polynomials

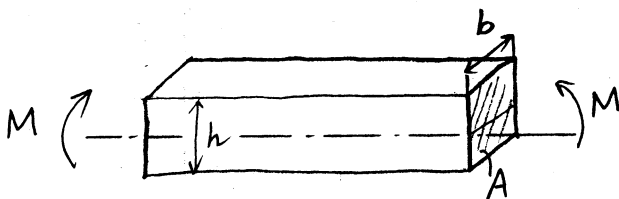
$$\phi(x, y) = \frac{1}{2} A x^2 + \frac{1}{2} B x y - C x y$$

$$\sigma_{xx} = \underline{\hspace{2cm}} \quad \sigma_{yy} = \underline{\hspace{2cm}} \quad \sigma_{xy} = \underline{\hspace{2cm}}$$



Write down the stress function for a rectangular bar under uniaxial tensile stress σ_0 .

$$\phi(x, y) = \underline{\hspace{2cm}}$$



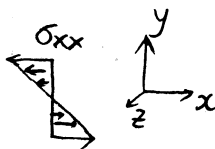
$$\sigma_{xx} \propto y$$

$$M = - \int_{-h/2}^{h/2} b \sigma_{xx}(y) \cdot y \, dy$$

$$\sigma_{xx} = - \frac{M y}{I_z}$$

$$I_z = \int_A y^2 dA = \frac{b h^3}{12}$$

↑ moment of inertia



Write down the stress function for a rectangular bar under pure bending M

$$\phi(x, y) = \underline{\hspace{2cm}}$$