

## Fundamental Equations of Elasticity:

compatibility, Equilibrium.

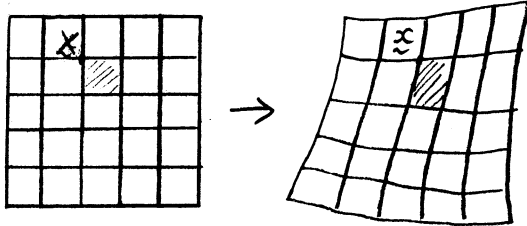
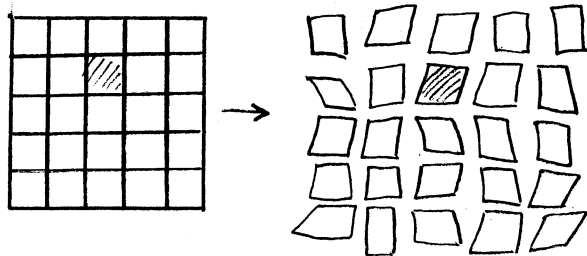
## §1. Compatibility condition for strains

 $u_i$  has 3 degrees of freedom (at every point  $\underline{x}$ ) $\epsilon_{ij}$  has 6 degrees of freedomgiven  $u_i(\underline{x})$ , we can always find  $\epsilon_{ij}(\underline{x})$  bydifferentiation:  $\epsilon_{ij}(\underline{x}) = \frac{1}{2}(u_{i,j} + u_{j,i})$ 

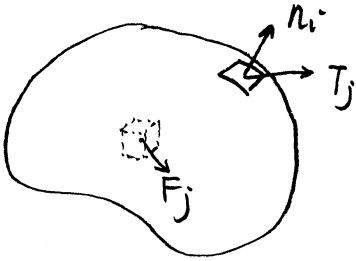
on the other hand, if we are given an arbitrary

 $\epsilon_{ij}(\underline{x})$ , we may not always be able to finda single valued, continuous  $u_j$ In order to be able to find a corresponding  $u_j(\underline{x})$  $\epsilon_{ij}(\underline{x})$  must satisfy some constraints.

compatibility condition: $\epsilon_{ij,k,l} + \epsilon_{kl,i,j} - \epsilon_{ik,j,l} - \epsilon_{jl,i,k} = 0$
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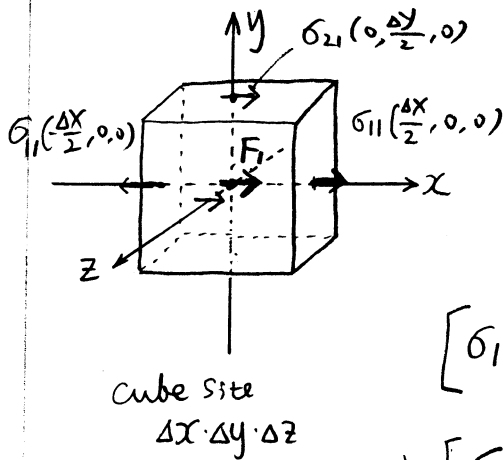
— This can be verified by plugging it into  $\epsilon_{ij} = \frac{1}{2}(u_{i,j} + u_{j,i})$ \* What is more difficult is to show this is a sufficient condition for us to find  $u_j$  corresponding to  $\epsilon_{ij}$ . (omitted here)Given a displacement field  $u_i(\underline{x})$ , it is straightforward to obtain the strain field.Given a strain field,  $\epsilon_{ij}(\underline{x})$ , we can imagine breaking the original medium into many pieces, each is deformed according to the local strain. — but the deformed pieces may not fit together (incompatible).

## §2. Equilibrium condition for stress



Consider a continuum body subjected to traction force  $T_j$  per unit area on the surface (normal vector  $n_i$ ) and body force  $F_j$  per unit volume

equilibrium condition:  $\sigma_{ij,i} + F_j = 0$



recall:  $\sigma_{ij}$  is force per unit area on  $i$ -th face of the stress cube in  $j$ -th direction

Total force balance in  $x$ -direction:

$$\begin{aligned} & \left[ \sigma_{11}\left(\frac{\Delta x}{2}, 0, 0\right) - \sigma_{11}\left(-\frac{\Delta x}{2}, 0, 0\right) \right] \cdot \underbrace{\Delta y \Delta z}_{\text{area}} \\ & + \left[ \sigma_{21}\left(0, \frac{\Delta y}{2}, 0\right) - \sigma_{21}\left(0, -\frac{\Delta y}{2}, 0\right) \right] \cdot \Delta x \Delta z \\ & + \left[ \sigma_{31}\left(0, 0, \frac{\Delta z}{2}\right) - \sigma_{31}\left(0, 0, -\frac{\Delta z}{2}\right) \right] \cdot \Delta x \Delta y \\ & + F_1(0, 0, 0) \cdot \Delta x \Delta y \cdot \Delta z = 0 \end{aligned}$$

In the limit of  $\Delta x, \Delta y, \Delta z \rightarrow 0$

$$\sigma_{11,1} + \sigma_{21,2} + \sigma_{31,3} + F_1 = 0$$

$$\sigma_{21,i} + F_1 = 0 \quad \text{— total force balance in } x\text{-direction.}$$

Similarly:  $\sigma_{i2,i} + F_2 = 0$  — total force balance in  $y$ -direction

$$\sigma_{i3,i} + F_3 = 0 \quad \text{— total force balance in } z\text{-direction.}$$

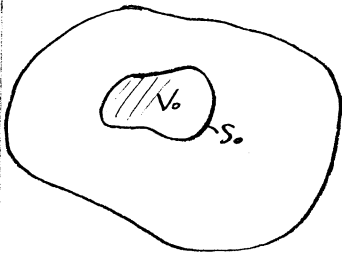
$$\therefore \sigma_{ij,i} + F_j = 0. \quad j \text{ is free index}$$

A more formal proof can be obtained using  
Gauss's Theorem (divergence theorem)



$$\int_S f_i n_i dS = \int_V f_{i,i} dV \quad (\text{Index notation})$$

$$\int_S \underline{f} \cdot \underline{n} dS = \int_V \nabla \cdot \underline{f} dV \quad (\text{vector notation})$$



consider an arbitrary volume  $V_0$   
 inside the elastic medium.

Force equilibrium requires

$$\underbrace{\int_{S_0} T_j dS}_{\text{traction force on the surface } S_0} + \underbrace{\int_{V_0} F_j dV}_{\text{body force inside } V_0} = 0$$

$$\int_{S_0} \sigma_{ij} n_i dS + \int_{V_0} F_j dV = 0$$

Apply Gauss's theorem, considering  $\sigma_{ij}$  as  $f_i$

$$\int_{V_0} (\sigma_{ij,i} + F_j) dV = 0$$

Because this is true for any volume  $V_0$ ,

$$\sigma_{ij,i} + F_j = 0$$

at every point inside the  
 continuum body.

In Summary, here are all the fundamental Equations of elasticity relating the following variables

displacement  $u_i$  — strain  $\epsilon_{ij}$  — stress  $\sigma_{ij}$  — traction force  $T_j$

definition of strain, stress:  $\epsilon_{ij} = \frac{1}{2} (u_{i,j} + u_{j,i})$

$$T_j = \sigma_{ij} n_i$$

compatibility:  $\epsilon_{ij,kl} + \epsilon_{kl,ij} - \epsilon_{ik,jl} - \epsilon_{jl,ik} = 0$

equilibrium:  $\sigma_{ij,i} + F_j = 0$

elastic constitutive law:

— Generalized Hooke's Law  
for isotropic medium

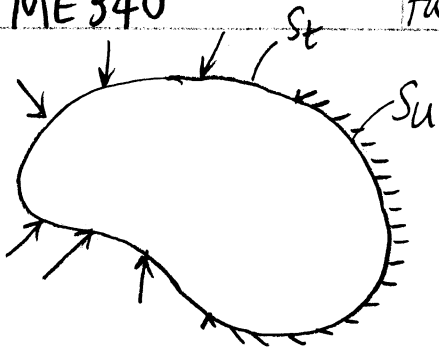
$$\sigma_{ij} = \lambda \epsilon_{kk} \delta_{ij} + 2\mu \epsilon_{ij}$$

$$\text{or } \epsilon_{ij} = \frac{1+\nu}{E} \sigma_{ij} - \frac{\nu}{E} \sigma_{kk} \delta_{ij}$$

All elasticity problems in this class satisfy the same equations as above.

Different applications correspond to different boundary conditions, which will lead to different solutions.

Boundary Value Problem (B.V.P.): Find the  $u_i, \epsilon_{ij}, \sigma_{ij}$  fields that satisfy compatibility, equilibrium, and constitutive equations and the specified boundary conditions.



In general, two types of boundary conditions can be applied on the surface.

traction boundary condition:  $T_j(\underline{x}) = q_j(\underline{x})$  on  $S_t$

i.e.  $\sigma_{ij}(\underline{x}) \cdot n_i(\underline{x}) = q_j(\underline{x})$  on  $S_t$

or

displacement boundary condition  $u_i(\underline{x}) = h_i(\underline{x})$  on  $S_u$ .

(or some combination of the two)

In the remaining of this class, we will

1. practice how to formulate the B.V.P.
2. learn the tools to solve the B.V.P.

General strategies to solve B.V.P.

1. Avoid compatibility eq. by working with  $u_i$  directly

(displacement formulation)

$$\sigma_{ij} = \lambda u_{k,k} \delta_{ij} + \mu (u_{i,j} + u_{j,i})$$

↓ plug into equilibrium cond.

$$\mu u_{i,kk} + (\lambda + \mu) u_{k,ki} + F_i = 0$$

$$\mu \nabla^2 \underline{u} + (\lambda + \mu) \nabla (\nabla \cdot \underline{u}) + \underline{F} = 0$$

Boundary Condition:

$S_u$ : is easy

$S_t$ : is a little more complicated.

mostly for 3D problems

2. stress formulation

need to work with compatibility condition.

rewrite compatibility cond. in terms of stress

+ equilibrium cond.  $\sigma_{ij,i} + F_j = 0$

⇒ Beltrami-Michell compatibility eq.

$$\sigma_{ij,kk} + \frac{1}{1+\nu} \sigma_{kk,ij} = -\frac{\nu}{1-\nu} \delta_{ij} F_{k,k} - F_{i,j} - F_{j,i}$$

mostly for 2D problems